## Chapter 36

\# 38. (a) Employing Eq. 36-3 with the small angle approximation ( $\sin \theta \approx \tan \theta=y / D$ where $y$ locates the minimum relative to the middle of the pattern), we find (with $m=1$ )

$$
D=\frac{y a}{m \lambda}=\frac{(0.90 \mathrm{~mm})(0.40 \mathrm{~mm})}{4.50 \times 10^{-4} \mathrm{~mm}}=800 \mathrm{~mm}=80 \mathrm{~cm}
$$

which places the screen 80 cm away from the slit.
(b) The above equation gives for the value of $y$ (for $m=3$ )

$$
y=\frac{(3) \lambda D}{a}=\frac{(3)\left(4.50 \times 10^{-4} \mathrm{~mm}\right)(800 \mathrm{~mm})}{(0.40 \mathrm{~mm})}=2.7 \mathrm{~mm} .
$$

Subtracting this from the first minimum position $y=0.90 \mathrm{~mm}$, we find the result $\Delta y=1.8 \mathrm{~mm}$.
\# 41. Consider two of the rays shown in Fig. 36-49, one just above the other. The extra distance traveled by the lower one may be found by drawing perpendiculars from where the top ray changes direction (point $P$ ) to the incident and diffracted paths of the lower one. Where these perpendiculars intersect the lower ray's paths are here referred to as points $A$ and $C$. Where the bottom ray changes direction is point $B$. We note that angle $\angle A P B$ is the same as $\psi$, and angle $B P C$ is the same as $\theta$ (see Fig. 36-49). The difference in path lengths between the two adjacent light rays is

$$
\Delta x=|A B|+|B C|=d \sin \psi+d \sin \theta .
$$

The condition for bright fringes to occur is therefore

$$
\Delta x=d(\sin \psi+\sin \theta)=m \lambda
$$

where $m=0,1,2, \ldots$ If we set $\psi=0$ then this reduces to Eq. 36-25.
\# 45. One strategy is to divide Eq. 36-25 by Eq. 36-3, assuming the same angle (a point we'll come back to, later) and the same light wavelength for both:

$$
\frac{m}{m^{\prime}}=\frac{m \lambda}{m^{\prime} \lambda}=\frac{d \sin \theta}{a \sin \theta}=\frac{d}{a} .
$$

We recall that $d$ is measured from middle of transparent strip to the middle of the next transparent strip, which in this particular setup means $d=2 a$. Thus, $m / m^{\prime}=2$, or $m=2 m^{\prime}$.

Now we interpret our result. First, the division of the equations is not valid when $m=0$ (which corresponds to $\theta=0$ ), so our remarks do not apply to the $m=0$ maximum. Second, Eq. 36-25 gives the "bright" interference results, and Eq. 36-3 gives the "dark" diffraction results (where the latter overrules the former in places where they coincide - see Figure 36-17 in the textbook). For $m^{\prime}=$ any nonzero integer, the relation $m=2 m^{\prime}$ implies that $m=$ any nonzero even integer. As mentioned above, these are occurring at the same angle, so the even integer interference maxima are eliminated by the diffraction minima.
\# 50. The condition for a minimum of a single-slit diffraction pattern is

$$
a \sin \theta=m \lambda
$$

where $a$ is the slit width, $\lambda$ is the wavelength, and $m$ is an integer. The angle $\theta$ is measured from the forward direction, so for the situation described in the problem, it is $1.20^{\circ}$ for $m=1$. Thus,

$$
a=\frac{m \lambda}{\sin \theta}=\frac{415 \times 10^{-9} \mathrm{~m}}{\sin 1.20^{\circ}}=1.98 \times 10^{-5} \mathrm{~m} .
$$

\# 51. The central diffraction envelope spans the range $-\theta_{1}<\theta<+\theta_{1}$ where $\theta_{1}=\sin ^{-1}(\lambda / a)$. The maxima in the double-slit pattern are located at

$$
\theta_{m}=\sin ^{-1} \frac{m \lambda}{d}
$$

so that our range specification becomes

$$
-\sin ^{-1}\left(\frac{\lambda}{a}\right)<\sin ^{-1}\left(\frac{m \lambda}{d}\right)<+\sin ^{-1}\left(\frac{\lambda}{a}\right)
$$

which we change (since sine is a monotonically increasing function in the fourth and first quadrants, where all these angles lie) to

$$
-\frac{\lambda}{a}<\frac{m \lambda}{d}<+\frac{\lambda}{a} .
$$

Rewriting this as $-d / a<m<+d / a$, we find $-6<m<+6$, or, since $m$ is an integer, $-5 \leq m \leq+5$. Thus, we find eleven values of $m$ that satisfy this requirement.

