Chapter 32

1. (a) A sketch of the field lines (due to the presence of the bar magnet) in the vicinity of the loop is shown below:



(b) The primary conclusion of Section 32-9 is two-fold: $\vec{\mu}$ is opposite to \vec{B} , and the effect of \vec{F} is to move the material toward regions of weaker magnetic fields. The direction of the magnetic moment vector (of our loop) is toward the right in our sketch, or in the +x direction.

(c) The direction of the current is clockwise (from the perspective of the bar magnet).

(d) Since the magnitude of *B* relates to the "crowdedness" of the field lines, we see that \vec{F} is toward the right in our sketch, or in the +*x* direction.

11. (a) Here, the enclosed electric flux is found by integrating

$$\Phi_E = \int_0^r E \, 2\pi r \, dr = t (0.500 \, \text{V/m} \cdot \text{s})(2\pi) \int_0^r \left(1 - \frac{r}{R}\right) r \, dr = t \pi \left(\frac{1}{2}r^2 - \frac{r^3}{3R}\right)$$

with SI units understood. Then (after taking the derivative with respect to time) Eq. 32-3 leads to (2 + 2)

$$B(2\pi r) = \varepsilon_0 \mu_0 \pi \left(\frac{1}{2}r^2 - \frac{r^3}{3R}\right).$$

For r = 0.0150 m and R = 0.0300 m, this gives $B = 2.78 \times 10^{-20}$ T.

(b) The integral shown above will no longer (since now r > R) have r as the upper limit; the upper limit is now R. Thus,

$$\Phi_{E} = t\pi \left(\frac{1}{2}R^{2} - \frac{R^{3}}{3R}\right) = \frac{1}{6}t\pi R^{2}.$$

Consequently, Eq. 32-3 becomes

$$B(2\pi r) = \frac{1}{6}\varepsilon_0\mu_0\pi R^2$$

which for r = 0.0400 m, yields

$$B = \frac{\varepsilon_0 \mu_0 R^2}{12r} = \frac{(8.85 \times 10^{-12})(4\pi \times 10^{-7})(0.030)^2}{12(0.0400)} = 2.09 \times 10^{-20} \text{ T}.$$

15. (a) Eq. 32-11 applies (though the last term is zero) but we must be careful with $i_{d,enc}$. It is the enclosed portion of the displacement current. Thus Eq. 32-17 (which derives from Eq. 32-11) becomes, with i_d replaced with $i_{d,enc}$,

$$B = \frac{\mu_0 i_{d \text{ enc}}}{2\pi r} = \frac{\mu_0 (5.00 \text{ A})(r/R)}{2\pi r}$$

which yields (after canceling *r*, and setting R = 0.0300 m) $B = 33.3 \mu$ T.

(b) Here
$$i_d = 3.00$$
 A, and we get $B = \frac{\mu_0 i_d}{2\pi r} = 16.7 \ \mu \text{T}$.

25. We use Eq. 32-27 to obtain

$$\Delta U = -\Delta(\mu_{s,z}B) = -B \ \Delta \mu_{s,z},$$

where $\mu_{s,z} = \pm eh/4\pi m_e = \pm \mu_B$ (see Eqs. 32-24 and 32-25). Thus,

$$\Delta U = -B\left[\mu_B - (-\mu_B)\right] = 2\mu_B B = 2\left(9.27 \times 10^{-24} \text{ J/T}\right)(0.48 \text{ T}) = 8.9 \times 10^{-24} \text{ J}.$$

26. (a) From Eq. 32-10,

$$i_{d} = \varepsilon_{0} \frac{d\Phi_{E}}{dt} = \varepsilon_{0} A \frac{dE}{dt} = \varepsilon_{0} A \frac{d}{dt} \Big[(9.0 \times 10^{5}) - (4.0 \times 10^{4} t) \Big] = -\varepsilon_{0} A (4.0 \times 10^{4} \text{ V/m} \cdot \text{s}) \\ = - (8.85 \times 10^{-12} \text{ C}^{2}/\text{N} \cdot \text{m}^{2}) (2.0 \times 10^{-2} \text{ m}^{2}) (4.0 \times 10^{4} \text{ V/m} \cdot \text{s}) \\ = -7.1 \times 10^{-9} \text{A}.$$

Thus, the magnitude of the displacement current is $|i_d| = 7.1 \times 10^{-9} \text{ A}$.

(b) The negative sign in i_d implies that the direction is downward.

(c) If one draws a counterclockwise circular loop *s* around the plates, then according to Eq. 32-18,

$$\sum_{s} \vec{B} \cdot d\vec{s} = \mu_0 i_d < 0,$$

which means that $\vec{B} \cdot d\vec{s} < 0$. Thus \vec{B} must be clockwise.

34. From Sample Problem 32.01 — "Magnetic field induced by changing electric field," we know that $B \propto r$ for $r \leq R$ and $B \propto r^{-1}$ for $r \geq R$. So the maximum value of *B* occurs at r = R, and there are two possible values of *r* at which the magnetic field is 60% of B_{max} . We denote these two values as r_1 and r_2 , where $r_1 < R$ and $r_2 > R$.

- (a) Inside the capacitor, $0.60B_{\text{max}}/B_{\text{max}} = r_1/R$, or $r_1 = 0.60R = 0.60(30 \text{ mm}) = 18 \text{ mm}$.
- (b) Outside the capacitor, $0.60B_{\text{max}}/B_{\text{max}} = (r_2/R)^{-1}$, or

$$r_2 = R/0.60 = 50$$
 mm.

(c) From Eqs. 32-15 and 32-17,

$$B_{\text{max}} = \frac{\mu_0 i_d}{2\pi R} = \frac{\mu_0 i}{2\pi R} = \frac{(4\pi \times 10^{-7} \text{ T} \cdot \text{m/A})(5.0 \text{ A})}{2\pi (0.040 \text{ m})} = 2.5 \times 10^{-5} \text{ T}.$$

36. From Eq. 28-11, we have $i = (\% / R) e^{-t/\tau}$ since we are ignoring the self-inductance of the capacitor. Equation 32-16 gives

$$B=\frac{\mu_0 i_d r}{2\pi R^2}\,.$$

Furthermore, Eq. 25-9 yields the capacitance

$$C = \frac{\varepsilon_0 \pi (0.05 \text{ m})^2}{0.003 \text{ m}} = 2.318 \times 10^{-11} \text{F},$$

so that the capacitive time constant is

$$\tau = (20.0 \times 10^6 \,\Omega)(2.318 \times 10^{-11} \,\mathrm{F}) = 4.636 \times 10^{-4} \,\mathrm{s}.$$

At $t = 750 \times 10^{-6}$ s, the current is

$$i = \frac{9.60 \text{ V}}{20.0 \times 10^6 \Omega} e^{-t/\tau} = 9.52 \times 10^{-8} \text{ A} .$$

Since $i = i_d$ (see Eq. 32-15) and r = 0.0300 m, then (with plate radius R = 0.0500 m) we find

$$B = \frac{\mu_0 i_d r}{2\pi R^2} = \frac{(4\pi \times 10^{-7} \,\mathrm{T \cdot m/A})(9.52 \times 10^{-8} \,\mathrm{A})(0.030 \,\mathrm{m})}{2\pi (0.0500 \,\mathrm{m})^2} = 2.28 \times 10^{-13} \,\mathrm{T}.$$